

Accurate, conformation-dependent predictions of solvent effects on protein ionization constants

P. Barth*[†], T. Alber*, and P. B. Harbury[‡]

*Department of Molecular and Cell Biology, University of California, Berkeley, CA 94720-3206; and [‡]Department of Biochemistry, Stanford University, Stanford, CA 94305

Communicated by Robert R. Baldwin, Stanford University Medical Center, Stanford, CA, January 9, 2007 (received for review July 12, 2006)

Predicting how aqueous solvent modulates the conformational transitions and influences the pKa values that regulate the biological functions of biomolecules remains an unsolved challenge. To address this problem, we developed FDPB.MF, a rotamer repacking method that exhaustively samples side chain conformational space and rigorously calculates multibody protein–solvent interactions. FDPB.MF predicts the effects on pKa values of various solvent exposures, large ionic strength variations, strong energetic couplings, structural reorganizations and sequence mutations. The method achieves high accuracy, with root mean square deviations within 0.3 pH unit of the experimental values measured for turkey ovomucoid third domain, hen lysozyme, *Bacillus circulans* xylanase, and human and *Escherichia coli* thioredoxins. FDPB.MF provides a faithful, quantitative assessment of electrostatic interactions in biological macromolecules.

conformational flexibility | electrostatics | pKa prediction | protein modeling

By strongly interacting with charges, the solvent significantly modulates the electrostatic properties of biomolecular systems. Although variations in pH and ionic strength are responsible for physiologically important conformational transitions (1), quantitatively assessing their role in modulating electrostatic energies is particularly challenging (2). Approaches in which solvent is modeled as a continuum polarizable medium while biomolecules are modeled in explicit molecular detail have become widely used in recent years (3, 4). However, no current method combines broad conformational sampling with a rigorous solvation model that can predict quantitatively and efficiently the effects of solvent on protein energetics and conformations.

In principle, molecular dynamics and free-energy simulations monitoring protonation events can model accurately the energetic effects of solvation and conformational relaxation of biomolecules. However, despite recent progress, these techniques often fail to converge in a reasonable amount of computer time and therefore do not provide reliable predictions of thermodynamic properties (5, 6).

Rotamer repacking methods sample more efficiently and exhaustively the conformational space of biomolecules. However, these methods require the energy function be decomposed into self energies and interaction energies between pairs of residues (7–10). Consequently, repacking methods cannot model the effects of conformational relaxations involving more than two positions at a time. Many important properties of biomolecular surfaces (i.e., their interface with the solvent and the distribution of bulk ions around them) are not pairwise factorable and depend on the simultaneous knowledge of the conformations of all residues. Approximating these effects by relaxing the protein–solvent interface only at the vicinity of solvent-exposed charged residues does not improve the prediction of pKa values in proteins (11). Current rotamer repacking methods have also difficulties in assessing accurately and reliably the electrostatic properties of buried charged residues that often play important roles in catalysis, structural specificity, and folding (2, 9–11).

To address these problems, we developed FDPB.MF, a rotamer repacking method that combines a general and full treatment of side chain conformational flexibility with the rigorous computation of multibody protein–solvent interactions. Non-pairwise factorable solvation energy terms calculated by a finite-difference Poisson–Boltzmann (FDPB) method were incorporated into a mean-field (MF) side chain rotamer repacking algorithm. Side chain conformations that minimize the free energy of the system were identified with the explicit relaxation of the protein–solvent interface and of the ion distribution around the protein. Here, we apply FDPB.MF to the classic problem of calculating the ionization constants of protein side chains.

Results

For further details, see [supporting information \(SI\) Text, SI Figs. 4–7, and SI Tables 5–8](#).

FDPB.MF was developed to provide a solution to the problem of combining efficiently the exhaustive sampling of protein side chain conformations with a rigorous treatment of solvation. The method is based on an ensemble description in which each conformation is assigned a weight corresponding to its probability of being occupied in the population. The effective quantities describing the probability-weighted conformational ensembles were derived so that they would equate or well approximate the corresponding average calculated by enumerating all of the discrete states of the system. During the development of the method, three important challenges were encountered that led to algorithmic developments for accurate calculation of the solvation energy of a probability-weighted conformational ensemble [Eq. 1 and [SI Fig. 5](#)], calculation of probability-dependent physical quantities to solve the Poisson–Boltzmann equation ([SI Text Eqs. S2 and S4](#)) and mapping of probability-weighted dielectric boundaries on a lattice ([SI Text Eqs. S6 and S7 and SI Fig. 6](#)).

Choice of the Solute Dielectric Constant. The choice of a dielectric constant for the solute (ϵ_p) depends on the level of explicit modeling of its polarizabilities. $\epsilon_p = 2$ is commonly assumed to be an adequate value for the implicit average treatment of electronic polarizabilities. Because electronic fluctuations and backbone conformational relaxations are not treated explicitly in FDPB.MF, the solute was assigned a higher ϵ_p of 4, a value derived from crystalline acetamide, a molecular analogue of the

Author contributions: P.B., T.A., and P.B.H. designed research; P.B. performed research; P.B., T.A., and P.B.H. analyzed data; and P.B. and T.A. wrote the paper.

The authors declare no conflict of interest.

Abbreviations: PB, Poisson–Boltzmann; FDPB, finite-difference PB; MF, mean-field; SCMF, self-consistent MF; OMTKY3, turkey ovomucoid third domain; BCX, *B. circulans* xylanase; TRX, *E. coli* thioredoxin; HEWL, hen lysozyme.

[†]To whom correspondence should be sent at the present address: Department of Biochemistry, University of Washington, Seattle, WA 98195. E-mail: barthp@u.washington.edu.

This article contains supporting information online at www.pnas.org/cgi/content/full/0700188104/DC1.

© 2007 by The National Academy of Sciences of the USA

Table 1. Effect of the solute dielectric constant on the pKa values calculated for OMTKY3

	Sites						rms error	Max error
	Asp-7	Glu-10	Glu-19	Asp-27	Glu-43	C terminus		
Experimental*	2.7 ± 0.06	4.1 ± 0.07	3.2 ± 0.02	2.2 ± 0.06	4.8 ± 0.1	<2.7 ± 0.12		
FDPB_MF† ($\epsilon_p^\ddagger = 4$)	2.5	3.9	3.0	2.3	4.4	2.7	0.2	0.4
FDPB_MF† ($\epsilon_p^\ddagger = 2$)	2.5	4.5	2.6	3.2	5.1	2.7	0.5	0.9

*Experimentally measured pKa values at 25°C and 10 mM monovalent salt (14).

†FDPB_MF, pKa values calculated with an ionic strength of 10 mM (see *Materials and Methods*).

‡ ϵ_p , the protein dielectric constant used for the calculation.

protein backbone (12). However, no sets of atomic charges and radii with $\epsilon_p = 4$ have been parameterized to reproduce solvation energies of peptides or proteins. The closest set of parameters is the PARSE set, parameterized to reproduce solvation energies of small molecules with solute polarizabilities corresponding to ϵ_p ranging from 2 to 3 (13). In principle, the calculations should be performed with the ϵ_p of the molecules from which the PARSE set of charges and radii were derived. Table 1 summarizes the prediction of the pKa values of OMTKY3 by FDPB_MF performed with $\epsilon_p = 2$ or 4. The predictions with $\epsilon_p = 2$ were still acceptable (rmsd = 0.5) but significantly less accurate than the ones performed with $\epsilon_p = 4$, especially for the buried Asp-27 (rms error of 0.9 versus 0.0). The main reasons for the reduced accuracy could be

1. FDPB_MF's treatment of polarizability due to conformational relaxation of the polypeptide chain is limited to the effects of sampling side chain rotameric degrees of freedom. The use of $\epsilon_p = 4$ may compensate for the lack of treatment of backbone conformational relaxation and covalent bond distortions and also for the discrete representation of the conformational space.
2. Solvation due to electronic polarizability may be higher in protein cores than in highly solvated small molecules and would explain why the loss of accuracy mainly affected the buried Asp-27.
3. With $\epsilon_p = 2$, the difference in solvation energies and electrostatic interactions between solvent-exposed and buried positions are higher than with $\epsilon_p = 4$. The ruggedness of the conformational free energy landscape may therefore be higher with $\epsilon_p = 2$ than with $\epsilon_p = 4$. Accordingly, the self-consistent MF (SCMF) algorithm did not always well converge at $\epsilon_p = 2$ and this may also explain why the predictions were less accurate. To ensure convergence, the lambda factor (SI Fig. 4 and ref. 15) which controls the rotamer probability updates in the SCMF was subsequently lowered to 0.15 in all of the calculations performed in this study. In these calculations, the solute and the solvent were also always assigned a dielectric constant of 4 and 80, respectively.

Test of the FDPB_MF Algorithm. To validate the physical basis of the approach, we performed energy calculations and conformational searches on a model OMTKY3 with three flexible polar sites that we could enumerate discretely. Effective solvation energies were found to be within 0.3 kcal/mol of the exact probability-weighted average energies of the discrete states of the system (data not shown). The ability of the FDPB_MF algorithm to find pH-dependent global free energy minima also was assessed. The FDPB_MF titration curves were compared with the exact curve obtained by exhaustively sampling all of the microstates of the system (Fig. 1). The pKa of Asp-27 predicted by FDPB_MF in the model OMTKY3 lies within 0.3 pH unit of the value calculated by using the exact scheme. However, the variations in the probability of the deprotonated species upon

pH predicted by FDPB_MF were much sharper than the one predicted by enumerating all of the microstates of the system. Asp-27 protonated and deprotonated species interact strongly with two different conformations of Tyr-31 (see SI Fig. 7A). This strong conformational coupling could not be resolved by a simple mean-field averaging of the conformational ensembles (zero- and first-order terms of Eq. 1) (16, 17). At pH values close to the pKa, where these mutually exclusive configurations have similar free energies, FDPB_MF only selects the configuration that has the lowest free energy. To accurately compute the probability of each chemical specie at pH values close to the pKa, protonated and deprotonated states for this residue pair have to be enumerated and treated separately (second-order term in Eq. 1). The second-order term in Eq. 1 also proved to be essential in the accurate prediction of the pKas of the two buried energetically coupled glutamate residues of *Bacillus circulans* xylanase.

Prediction of Solvent-Dependent Protein Ionization Constants. We then assessed the ability of FDPB_MF to accurately predict solvent-dependent energetics in proteins. For all FDPB_MF predictions, conformational flexibility was modeled for side chains atoms, whereas backbone atoms were fixed to the coordinates determined by x-ray crystallography. The resulting conformational ensemble was relaxed by the SCMF algorithm to find the distribution of rotamer probabilities that minimizes the free energy of the system. The pKa values were assigned to the pH at which the corresponding sites become 50% deprotonated.

Fig. 2 summarizes the pKa calculations performed by FDPB_MF on 31 acidic residues and compares these values to the predictions published by other methods. Among these,

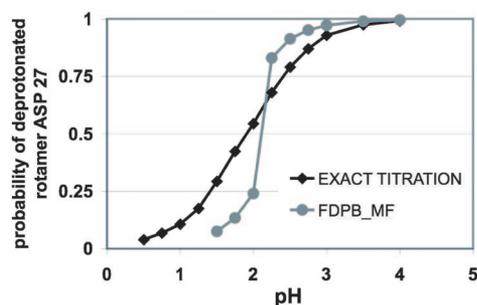


Fig. 1. Accuracy of the convergence of the FDPB_MF compared with the solution obtained by exhaustively sampling all of the discrete states of the system. The two different predicted titration curves of Asp 27 correspond to the sum of the probability of the deprotonated rotamers of Asp-27 as a function of the pH. The "exact titration" curve (black) was obtained by enumerating all of the discrete states of the system and is equal to the sum at a given pH of the Boltzmann weights of the discrete states occupied by the deprotonated rotamers. The FDPB_MF curve (gray) is derived by relaxing a rotameric probability distribution with the FDPB_MF method. It corresponds to the probability that minimizes at a given pH the electrostatic free energy of the protein after optimization by the SCMF algorithm.

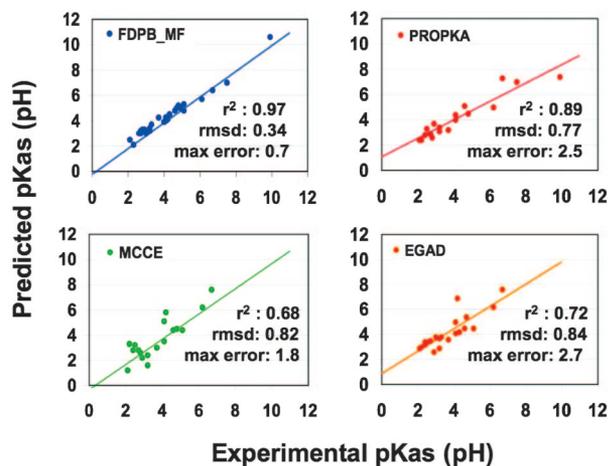


Fig. 2. Synopsis of the pKa predictions performed by FDPB.MF and comparison with other methods: PROPKA (18), MCCE (10), and EGAD (9). rmsds and maximal errors to the experimental values are in pH units. Each line and r^2 value corresponds to the best linear regression fit to the data and the correlation coefficient, respectively.

PROPKA performs best, with a rmsd, maximal error to the experimental values and correlation coefficient of 0.77 pH unit, 2.5 pH units, and 0.89, respectively. However, the slope of the best fit to these predictions was significantly lower than 1.0 (0.73), suggesting that the model underlying PROPKA has been trained to reproduce experimental data rather than accurately recapitulating first physical principles. FDPB.MF predicts the pKa values with a rmsd, maximal error to the experimental

values and correlation coefficient of 0.34 pH unit, 0.70 pH unit, and 0.97, respectively. The slope of the best fit to these predictions was 1.01. We largely focused our study on functionally or structurally important residues with pKas that are often difficult to predict accurately (Table 2). PROPKA surpasses other published methods with a rmsd and maximal error to the experimental values of 0.89 unit and 2.5 pH units, respectively. FDPB.MF predicts the pKa values with a rmsd, maximal error to the experimental values of 0.38 and 0.70 pH unit, respectively.

The ionic strength is an important parameter governing the stability and function of proteins. Several acidic sites of OMTKY3 show measured pKa shifts of 0.25–0.77 pH unit upon variation of the ionic strength from 10 mM to 1 M (19). FDPB.MF accurately predicted the effects of such large variation in ionic strength (Table 3). The rmsd and the maximum error to the experimental values were 0.14 and 0.2 pH unit, respectively. The sensitivity of pKas to the change in bulk ionic strength often reflects the nature of the electrostatic interactions stabilizing a particular ionized site (30, 31). These results suggest that FDPB.MF predicts quantitatively the complex balance of electrostatic interactions that stabilize charged residues.

Accurately predicting the energetic effects of mutations is an important goal in protein modeling. The insensitivity of the pKas of three acidic residues (Asp-7, Glu-10, Glu-19) to the removal of a neighboring positively charged residue (Lys 34) was not accurately predicted by simple FDPB calculations (20). FDPB.MF predicted smaller pKa perturbations (Table 4), with rmsd and maximal error from the experimental data of 0.28 and 0.4 pH unit, respectively.

Electrostatic interactions play a crucial role in modulating the catalytic reactivity of enzymes. We assessed the ability of FDPB.MF to accurately predict the pKa values of the two buried energetically coupled glutamate residues that lie at the heart of

Table 2. Comparison between experimental and calculated pKa values for structurally and functionally important sites

Protein	Sites	Experiment*	FDPB.MF [†] ($\epsilon_p^{\ddagger}=4$)	PROPKA [§]	MCCE ($\epsilon_p^{\ddagger}=4$)	EGAD ($\epsilon_p^{\ddagger}=8$)	FDPB** ($\epsilon_p^{\ddagger}=20$)	Null model ^{††}
OMTKY3 <i>I</i> = 10 mM	D7	2.7	3.1	2.8	2.8	3.1	2.9	4.0
	D27	2.2	2.1	2.4	3.3	3.0	3.6	4.0
	E19	3.2	3.2	3.1	1.6	3.7	2.6	4.4
<i>I</i> = 1 M	D7	3.3	3.7	ND	ND	ND	3.5	4.0
	D27	2.9	3.0	ND	ND	ND	4.0	4.0
	E19	4.0	3.9	ND	ND	ND	3.0	4.4
HEWL	D52	3.7	4.3	3.2	3.0	3.6	3.1	4.0
	D87	2.1	2.5	2.4	1.2	2.9	2.7	4.0
	E7	2.9	3.3	3.7	2.2	2.6	3.3	4.4
BCX	E35	6.2	5.7	5.0	6.2	6.2	4.4	4.4
	E78	4.6	4.8	5.1	ND	4.4	2.9	4.4
	E172	6.7	6.4	7.3	ND	7.6	5.9	4.4
E172Q	E78	5.1	5.0	ND	ND	4.5	4.3	4.4
E78Q	E172	4.2	3.8	ND	ND	6.9	2.8	4.4
TRX	D26	7.5	7.0	7.0	ND	ND	5.8	4.0
TRH	D26	9.9	10.6	7.4	ND	ND	7.5	4.0
rmsd			0.38	0.89	0.92	0.98	1.24	2.11
Max error			0.7	2.5	2.6	2.7	2.4	5.9

ND, not determined.

*Experimentally measured pKa values (14, 19, 20–24).

[†]FDPB.MF, pKa values calculated with FDPB.MF (see *Materials and Methods*).

[‡] ϵ_p , the protein dielectric constant used for the calculation.

[§]PROPKA, pKa values calculated with an empirical structure-based pKa prediction method (18).

^{||}MCCE, pKa values calculated with a pairwise factorable FDPB electrostatic potential and with polar side-chain rotamer conformations relaxed by a Monte Carlo algorithm (10).

^{||}EGAD, pKa values calculated with a pairwise factorable Generalized-Born electrostatic potential and with side-chain rotamer conformations relaxed by a SCMF algorithm (9).

**FDPB, pKa values calculated by using an FDPB electrostatic potential of mean force and discrete backbone and side-chain coordinates from the x-ray crystal structures (25–29).

^{††}Intrinsic model compound pKa values (8).

Table 3. Experimental and calculated ionic strength effects on the pKa values of OMTKY3

	Sites					rms error	Max error
	Asp-7	Glu-10	Glu-19	Asp-27	Glu-43		
Experimental*	0.57 ± 0.2	0.25 ± 0.14	0.77 ± 0.05	0.72 ± 0.1	0.30 ± 0.14		
FDPB_MF† (ε _p = 4)	0.63	0.40	0.74	0.90	0.10	0.14	0.25
FDPB‡ (ε _p = 20)	0.58	0.54	0.41	0.43	0.14	0.21	0.47

*Differences between the experimentally measured pKa values at 25°C and 1 M monovalent salt (19) and the same values measured at 10 mM monovalent salt.

†FDPB_MF, differences between the pKa values computed with an ionic strength of 1 M bulk and the same values computed with an ionic strength of 10 mM.

‡FDPB, literature pKa shifts calculated by using an FDPB electrostatic potential of mean force (19).

the active site of *B. circulans* xylanase (BCX; Table 2 and SI Table 8). Simple FDPB calculations with a single dielectric constant of 8 for the protein failed to predict accurately the pKa values of both glutamates in the WT protein. By introducing an explicit treatment of conformational flexibility, EGAD performed better than FDPB and similar to PROPKA for the WT protein. However, EGAD failed to predict the large pKa shift (−2.5 pH units) observed for Glu-172 when Glu-78 was mutated to Gln (21, 22). By finely sampling the conformational space in the vicinity of the glutamates and by uncoupling their titration in the WT protein, FDPB_MF accurately predicted the unusually high pKa of Glu-172 (pKa of 6.7) and its large pKa shift upon mutation of Glu-78 to Gln. The rmsd and maximal error to the experimental values were equal to 0.2 and 0.3 pH unit, respectively.

The buried, catalytically important Asp-26 in *Escherichia coli* and human thioredoxins have among the highest pKa values ever measured for an aspartate (23). FDPB_MF predicted their pKa more accurately than other methods, with errors of 0.5 and 0.7 pH unit compared with the experimental values (Table 2 and SI Table 8).

Prediction of Solvent-Dependent Protein Conformations. Because protein energetics and conformations are tightly coupled, we expect quantitative energetic predictions to be corroborated with high-resolution structural predictions.

Fig. 3 compares the observed and predicted structural reorganizations occurring in the active site of *B. circulans* xylanase upon titration of Glu 172. Although the crystals were soaked at pH 4.0, the observed conformation at “pH 4.0” likely corresponds to that induced by deprotonated Glu 78 and protonated Glu 172 (21) and can be compared with the structure predicted at pH 5.5. Except for a small difference at Tyr 69, the observed and predicted conformational changes induced by protonation of Glu 172 are in good agreement. The aromatic ring of Tyr 80 flips by 180° to preferentially stabilize the remaining deprotonated

Glu at position 78. The Asn 35 amide group moves away from the protonated Glu 172 and the amide of Gln 127 comes slightly closer to Glu 78. The rmsd between the predicted and observed distance changes involving protons and heteroatoms in the active site was 0.2 Å (SI Table 5).

FDPB_MF also predicted substantial pH dependent conformational shifts (SI Text and SI Fig. 7) that have yet to be tested by experiments. Most of these conformational changes were corroborated with significantly improved energetic predictions when compared with traditional FDPB calculations performed on fixed crystallographic structures (Tables 2–4 and SI Tables 6–8).

Discussion

Protein stability, solubility, recognition, and catalysis depend critically on electrostatic interactions, and many important conformational transitions in proteins are triggered physiologically by changes in bulk ion composition and in pH. FDPB_MF was developed to accurately predict this coupling between solvent-dependent protein energetics and conformations. Here we used FDPB_MF to predict the pKa values of 31 carboxylates in very different protein environments and solvent conditions. This approach yielded predictions of unprecedented accuracy as well as new capabilities to model the effects of mutations and changes in ionic strength. The accuracy of the pKa predictions approached the experimental error in the measured values, adding considerable utility to the pKa predictions based on structural data.

For charged residues near the protein surface, the solvent accessibility and solvation energies are mainly determined by the conformation of neighboring residues or substrates. By explicitly modeling the conformational relaxations of all residues and/or ligands that constitute the protein-solvent interface, the FDPB_MF method captures the energetics of partially solvent-exposed acidic residues more accurately than available methods that rely solely on pairwise-factorable energy functions. Buried charged residues are less frequent and are often involved in

Table 4. Experimental and calculated effects of point mutations on the pKa values of OMTKY3

Mutant	Sites				rms error	Max error
	Asp-7	Glu-10	Glu-19			
Lys34Gln						
Experimental*	−0.1 ± 0.1	0.15 ± 0.14	0.1 ± 0.1			
FDPB_MF† (ε _p = 4)	0.3	0.4	0.2	0.28	0.4	
FDPB‡ (Xray, ε _p = 20)	0.6	0.4	0.6	0.52	0.7	
Lys34Thr						
Experimental*	0 ± 0.3	0.15 ± 0.1	−0.1 ± 0.1			
FDPB_MF† (ε _p = 4)	0.3	0.4	0.2	0.28	0.3	
FDPB‡ (X-ray, ε _p = 20)	0.6	0.4	0.6	0.55	0.7	

*Differences between the experimentally measured pKa values of the mutant protein and the same values measured for the WT protein (20).

†FDPB_MF: differences between the pKa values computed for the mutant protein and the same values computed for the WT protein.

‡FDPB: literature pKa shifts calculated using an FDPB electrostatic potential of mean force (20).

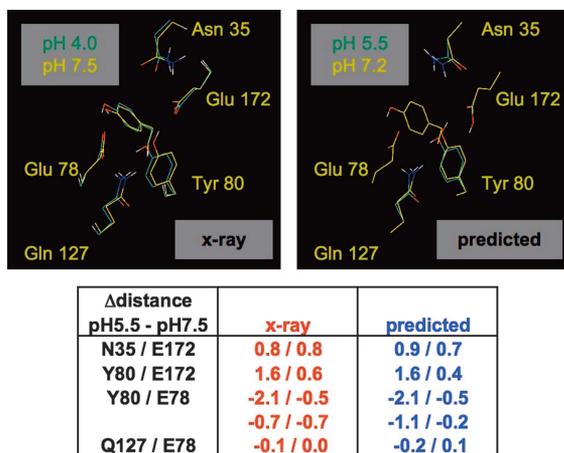


Fig. 3. Good agreement between observed and predicted conformational changes induced by the protonation of Glu 172 in *B. circulans* xylanase. Figures were generated by using Pymol (www.delanoscientific.com). As discussed in ref. 21, the crystal structure solved at an apparent pH of 4.0 is likely to be consistent with Glu 172 being protonated. Protons on the x-ray structures were added with the program Reduce (32). At the bottom, distance changes are provided in angstroms. The first and second numbers represent distances between protons and heteroatoms and distances between heteroatoms, respectively.

catalysis or in structural specificity. Anisotropy in polarizability of protein interiors is higher than at the surfaces, in line with the inability of FDPB calculations to capture implicitly the structural relaxation of protein interiors with a single dielectric constant (21). The energetics of buried residues are therefore sensitive to the fine conformational reorganization of their packed micro-environments. FDPB.MF explicitly models these effects and predicts the energetics of buried charged residues more accurately than current molecular dynamics simulations and rotamer repacking methods. The accuracy of FDPB.MF arises from the unique combination of fine sampling of the side chain conformational space, rigorous calculation of multibody protein-solvent interactions and a general treatment of conformational relaxation. Unlike current methods based on first principles, FDPB.MF quantitatively predicts the energetic effects of various protein environments, sequence mutations and solvent compositions (Fig. 2, rmsd from experimental values = 0.34 pH unit). Each calculation reported here took 2–6 days on a single Pentium 4 2.4 GHz CPU for a single pH.

Understanding the relationships between protein energetics, conformations and dynamics also presents major challenges in protein modeling. Proteins exhibit side chain conformational flexibility that is dictated by their structural environment and by the physical properties of the solvent. FDPB.MF predicts pH-induced, side chain conformational changes involving either changes in the dominant rotamers or redistribution of probabilities within a conformational ensemble (Fig. 3 and SI Fig. 7). These shifts highlight the complex nature of the electrostatic interactions involved in stabilizing charged residues. The explicit modeling of these structural reorganizations by FDPB.MF is corroborated by improved energetic predictions when compared with fixed-structure calculations (Tables 2–4 and SI Tables 6–8). The pH-induced conformational changes predicted by FDPB.MF in *B. circulans* xylanase are in good agreement with those observed in the crystals (Fig. 3). As the challenges of protein structure prediction continue to move toward the generation of high-resolution models (33), the accurate prediction of solvent-dependent conformational changes will grow in importance.

Materials and Methods

FDPB Calculations. Perturbation theory is an effective approach for the computation of mean-field properties and correlations

beyond the mean-field. Following the same formalism (i.e., in the form of a perturbation series), Eq. 1 describes a general solution to the electrostatic protein-solvent interaction energy (E^{FDPB}) of a conformational ensemble defined by a rotamer probability distribution ($P_{I,R}$).

$$E^{FDPB} = E^{DIST} + \sum_I \sum_R P_{I,R} (E_{I,R}^{100} - E^{DIST}) + \sum_I \sum_{J \neq I} \sum_R \sum_{R'} P_{I,R} P_{J,R'} (E_{IR, JR'}^{100} - (E_{IR}^{100} + E_{JR'}^{100})) + \dots \quad [1]$$

In this ensemble description, each rotamer is assigned a weight corresponding to its probability of being occupied in the population of rotamers. In the mean-field approximation, individual side chain rotamers interact with the probability-weighted average of all of the rotamers at the neighboring sites. The corresponding mean-field electrostatic protein-solvent interaction energies are defined by the zero- and first-order terms of Eq. 1. These were computed in two steps by the FDPB module (SI Fig. 5). First, all side chain rotamers R at all sites I weighted by their probabilities $P_{I,R}$ were mapped onto a three-dimensional grid. The Poisson-Boltzmann (PB) equation was solved by finite-differences for this “distributed” conformational ensemble and a corresponding solvation energy, E^{DIST} , was computed (zero-order term in Eq. 1). This “distributed” representation of the system is inaccurate partly because rotamers from the same site are present and interact with each other. An additional energy term for each rotamer, E_{IR}^{100} (first-order term in Eq. 1), is necessary and was computed by iteratively placing single discrete rotamers at each site and solving the PB equation for the solvation energy. Adding the first-order to the zero-order term of Eq. 1 corrects for the simultaneous presence of rotamers on the grid and gives an accurate solution to the mean-field protein-solvent electrostatic interaction energies. The mean-field treatment was sufficient for most calculations performed in this study. It failed, however, when side chain conformations from different sites were strongly coupled. Higher-order terms in Eq. 1 are then necessary to describe the mutually exclusive combinations of rotamers.

The mean-field treatment requires the definition of effective physical quantities that accurately describe the probability-weighted conformational ensembles. These quantities were derived so that they would equate or well approximate the corresponding average calculated by enumerating all of the discrete states of the system. If a rotamer R alone occupies a particular space with a probability P_R , it is equivalent to that space being occupied a fraction P_R of the time by the solute and a fraction $(1 - P_R)$ of the time by the solvent. Consequently, if E_R , E_S , and E_W define the Born solvation energy of a charge covered by rotamer R , pure solute, or water, respectively, then E_R should equate the probability weighted average of E_S and E_W as defined by Eq. 2

$$E_R = P_R E_S + (1 - P_R) E_W. \quad [2]$$

Eq. 2 was used to derive the effective physical quantities needed to solve the PB equation by finite differences for the probability-weighted conformational ensembles (SI Text).

The FDPB algorithm is based on the QNIFFT program (academic version of Delphi provided by K. Sharp, University of Pennsylvania, Philadelphia, PA). All calculations were carried out on a 129 cubic grid. Focusing yielded grid resolutions > 2.4 grid units/Å (34). The linearized form of the PB equation was solved by the inexact Newton method with a multilevel solver algorithm to aid convergence (35). The solute was mapped on

